



# Search for Violation of CPT and Lorentz invariance in $B_s$ meson oscillations.

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## Why Search for CPT Violation?

- Test the CPT theorem:
  - ... any Lorentz invariant local quantum field theory with a Hermitian Hamiltonian must have CPT symmetry...
- Standard Model Extension (SME) provides a framework for potential Lorentz and CPT invariance violation
  - Phys. Rev. D 39, 683 (1989), Nucl. Phys. B 359, 545 (1991)
  - Violations occur at Plank scale whilst still producing observable effects
- Searches already carried out in other neutral meson systems:  $K^0$  (Kloe),  $D^0$  (Focus),  $B_d^0$  (Belle and BaBar)
- Charge asymmetries in single and like-sign dimuon events could explained by CPT violation.  
V. Kostelecký and R Van Kooten Phys. Rev. D82: 101702, 2010

# Neutral Meson Mixing

CPT & Lorentz violating:  
involves differences between  
*diagonal* terms

$$\mathbf{H} = \begin{pmatrix} \langle M^0 | H | M^0 \rangle & \langle M^0 | H | \bar{M}^0 \rangle \\ \langle \bar{M}^0 | H | M^0 \rangle & \langle \bar{M}^0 | H | \bar{M}^0 \rangle \end{pmatrix} = \begin{pmatrix} M_{11} - \frac{i}{2}\Gamma_{11} & M_{12} - \frac{i}{2}\Gamma_{12} \\ M_{12}^* - \frac{i}{2}\Gamma_{12}^* & M_{22} - \frac{i}{2}\Gamma_{22} \end{pmatrix}$$

Eigenvalue equation for  $\mathbf{H}$  yields two eigenstates

- well defined masses and decay widths (lifetimes) correspond to physical states

If we have CPT violation the eigenstates can be written as

$$|B_{sL}\rangle \propto p\sqrt{1 - \xi_s}|B_s^0\rangle + q\sqrt{1 + \xi_s}|\bar{B}_s^0\rangle,$$

$$|B_{sH}\rangle \propto p\sqrt{1 + \xi_s}|B_s^0\rangle - q\sqrt{1 - \xi_s}|\bar{B}_s^0\rangle.$$

where  $\xi$  is zero if CPT is conserved and is given by.

$$\xi_s = \frac{(M_{11} - M_{22}) - \frac{i}{2}(\Gamma_{11} - \Gamma_{22})}{-\Delta m_s + \frac{i}{2}\Delta\Gamma_s/2} \approx \frac{\beta^\mu \Delta a_\mu}{-\Delta m_s + i\Delta\Gamma_s/2},$$

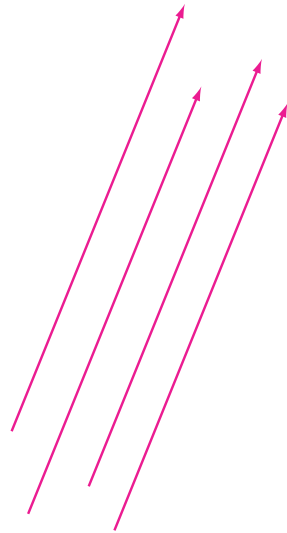
**tiny  $\rightarrow$  sensitivity**

$$\Delta m_s = 17.69 \pm 0.08 \text{ ps}^{-1} \quad 1 \text{ sec} = 1.52 \times 10^{24} \text{ GeV} \quad \Delta m_s = 1.16 \times 10^{-2} \text{ eV}$$



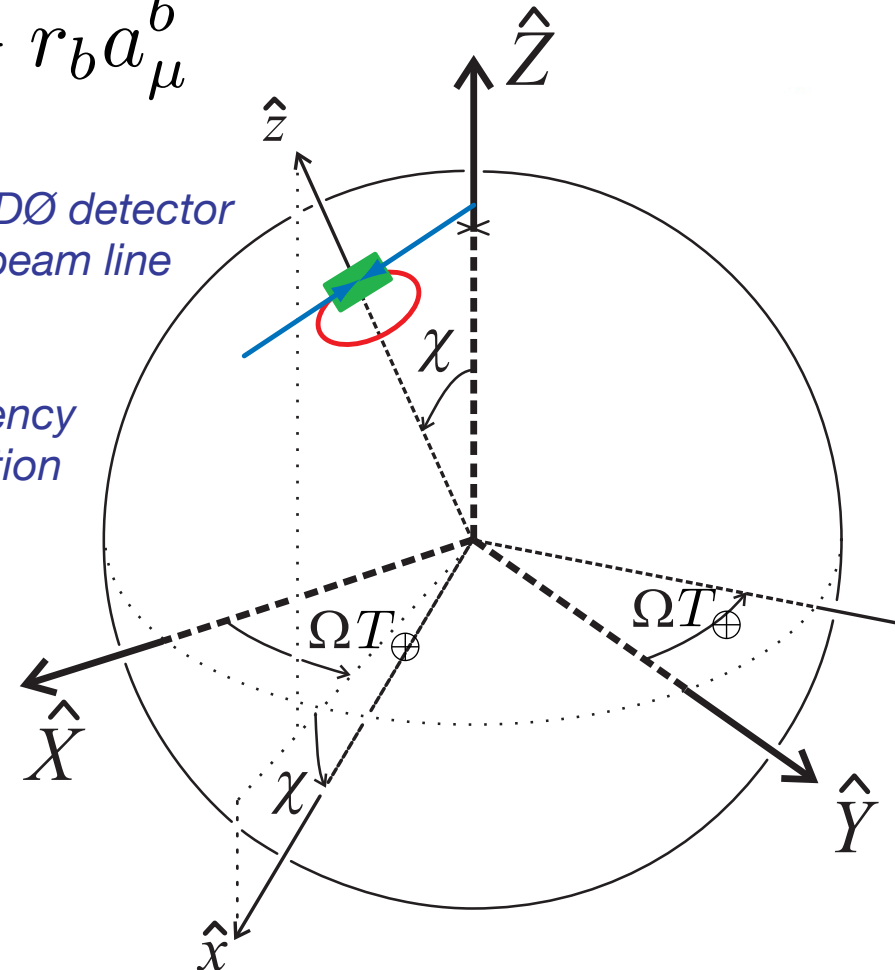
# Sidereal Time Variation

- Size of effect:  $\Delta a_\mu = r_s a_\mu^s - r_b a_\mu^b$



Constant vector field  $a_\mu$  in universe

Orientation of DØ detector and Tevatron beam line changes w.r.t. vector field  $a_\mu$  with the frequency of Earth's rotation



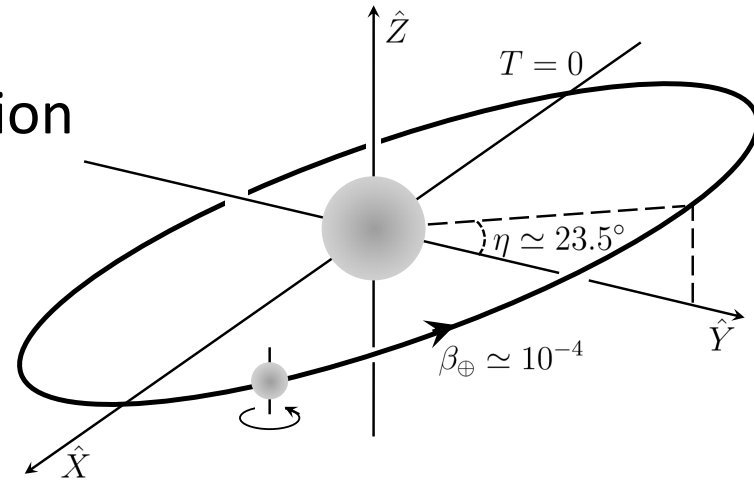
$$\Omega = 2\pi / (23^h 56^m 04.0982^s) \text{ Earth's sidereal frequency}$$

$T_\oplus$  local sidereal time of the collision event (time stamp)



# Experimental Effect

- We choose (T,X,Y,Z) as coordinates in the Sun-centered frame with the rotation axis of the Earth the Z-axis and X(Y) is at right ascension 0° (90°).
- T = 0: vernal equinox of the year 2000
- Transform to local (D0) coordinates



$$\mathcal{A}_{\text{CPT}} = \frac{-\Delta\Gamma_s \gamma^{\text{D0}}}{\Gamma_s \Delta m_s} \left[ \Delta a_T - C_\alpha S_\chi \beta_z^{\text{D0}} \Delta a_Z + \sqrt{C_\alpha^2 C_\chi^2 + S_\alpha^2} \sin(\Omega \hat{t} + \delta + \kappa) \beta_z^{\text{D0}} \Delta a_\perp \right],$$



# Experimental Effect

$$\mathcal{A}_{\text{CPT}} = \frac{-\Delta\Gamma_s \gamma^{\text{D}0}}{\Gamma_s \Delta m_s} \left[ \Delta a_T - C_\alpha S_\chi \beta_z^{\text{D}0} \Delta a_Z \right. \\ \left. + \sqrt{C_\alpha^2 C_\chi^2 + S_\alpha^2} \sin(\Omega \hat{t} + \delta + \kappa) \beta_z^{\text{D}0} \Delta a_\perp \right],$$

Looking for

$$\mathcal{A}_{\text{CPT}} = \text{Constant} + \text{Sin}(t + c)$$

where:  $\Delta a_\perp = \sqrt{\Delta a_X^2 + \Delta a_Y^2}$   $\delta = \text{atan}(\Delta a_X / \Delta a_Y)$

$$\kappa = \text{atan2}(C_\alpha C_\chi, -S_\alpha)$$

$$C_x = \cos(x), S_x = \sin(x)$$

$\chi$  = colatitude of D0 detector (angle from north pole)

$\alpha$  = proton beam direction

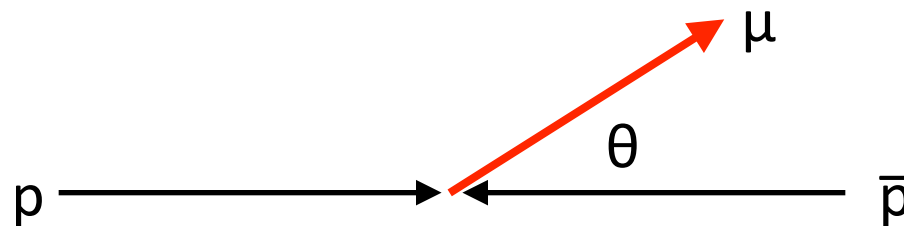


# $B_s^0$ Oscillations

- Using the decay process  $B_s \rightarrow D_s^- \mu^+ X$  in the analysis described in Phys. Rev. Lett., 110, 011801 (2013).
- Unlike other charge violation measurements looking at particle to particle oscillations:  $B_s \rightarrow B_s$
- Measure a raw asymmetry in bins of sidereal phase

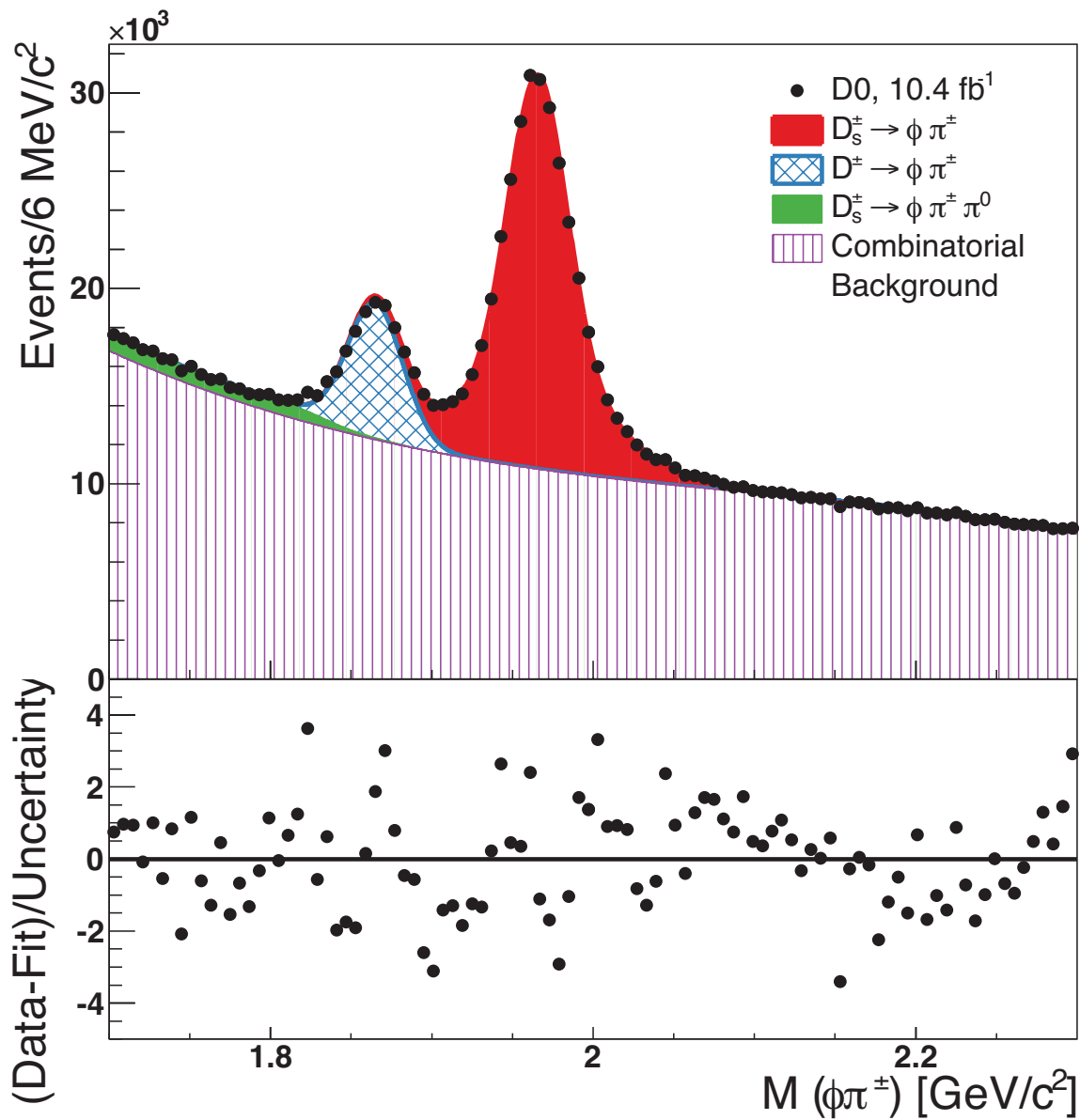
$$A = \frac{N_+ - N_-}{N_+ + N_-}$$

where  $N_{+(-)}$  is the number of events where  $Q^\mu \text{sgn}(\cos \theta) > 0$





# Dataset



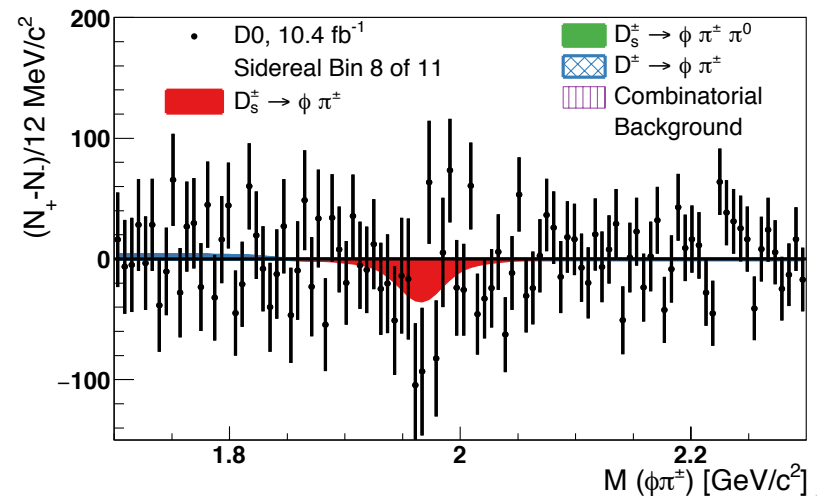
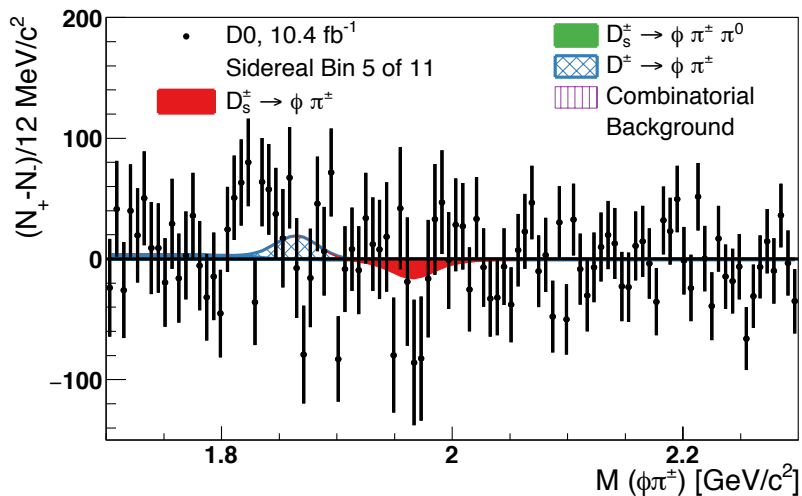




# Sensitivity



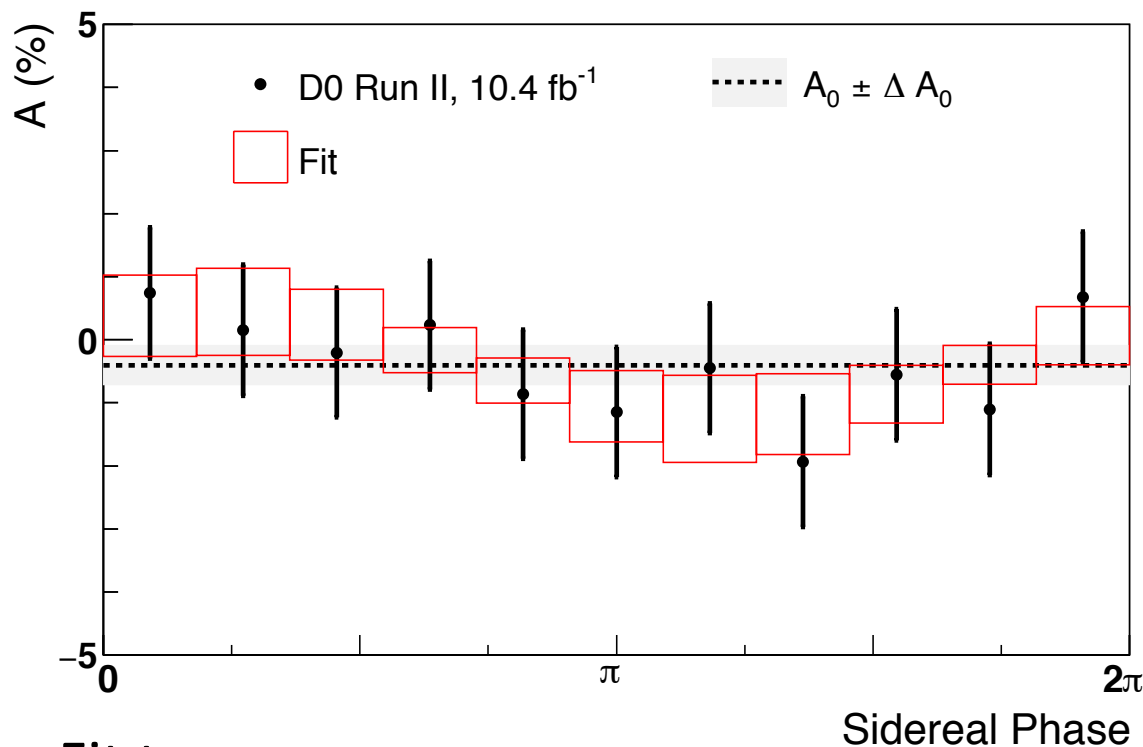
- Use charge injection to optimise number of bins in sidereal phase (using asymmetries of 0, 1, 2 and 5%) - Use 11 bins
- Fit all sidereal bins simultaneously
  - Common Signal peaks and background
  - Only Asymmetry of D<sub>s</sub> peak and D peak varies as function of sidereal phase.





# Sidereal Bins

$A_{Ds}$		
0.74%	$\pm$	1.03%
0.15%	$\pm$	1.03%
-0.20%	$\pm$	1.02%
0.23%	$\pm$	1.01%
-0.86%	$\pm$	1.02%
-1.14%	$\pm$	1.02%
-0.45%	$\pm$	1.02%
-1.93%	$\pm$	1.03%
-0.55%	$\pm$	1.03%
-1.11%	$\pm$	1.03%
0.68%	$\pm$	1.03%



Fit to:

$$A(\hat{t}) = A_0 - A_1 \sin(\Omega \hat{t} + \phi),$$

Significance (Wilk's Theorem):  
0.58 $\sigma$  deviation from zero (constant).



# Limit Setting

- Convert  $A_1$  into limit on  $a_{\perp}$

$$A_1 \sin(\Omega \hat{t} + \phi) = \frac{F_{B_s^0}^{\text{non-osc}} \Delta \Gamma_s \langle \gamma^{D^0} \beta_z^{D^0} \rangle}{\Gamma_s \Delta m_s} \times \sqrt{C_{\alpha}^2 C_{\chi}^2 + S_{\alpha}^2} \sin(\Omega \hat{t} + \delta + \kappa) \Delta a_{\perp},$$

where  $F^{\text{non-osc}}$  is the fraction of  $D_s^{\pm} \rightarrow \phi \pi^{\pm}$  decays for which an observed  $B_s^0$  has the same flavour as at birth

- giving

$$\Delta a_{\perp} < 1.2 \times 10^{-12} \text{ GeV}$$

- Convert  $A_0$  into a limit on  $a_T$  and  $a_Z$

$$(-0.8 < \Delta a_T - 0.396 \Delta a_Z < 3.9) \times 10^{-13} \text{ GeV}$$



# Comparison with Dimuon

$$(-0.8 < \Delta a_T - 0.396 \Delta a_Z < 3.9) \times 10^{-13} \text{ GeV}$$

- V.A. Kostelecky' and R.J. Van Kooten, Phys. Rev. D. 82, 101702(R) (2010), arXiv:1007.5312 makes a prediction of this value to explain the dimuon asymmetry:

$$(\Delta a_T - 0.396 \Delta a_Z) \simeq \times (3.7 \pm 3.8) 10^{-12} \text{ GeV}$$

so need effect of order  $10^{-12}$  GeV to make significant contribution.

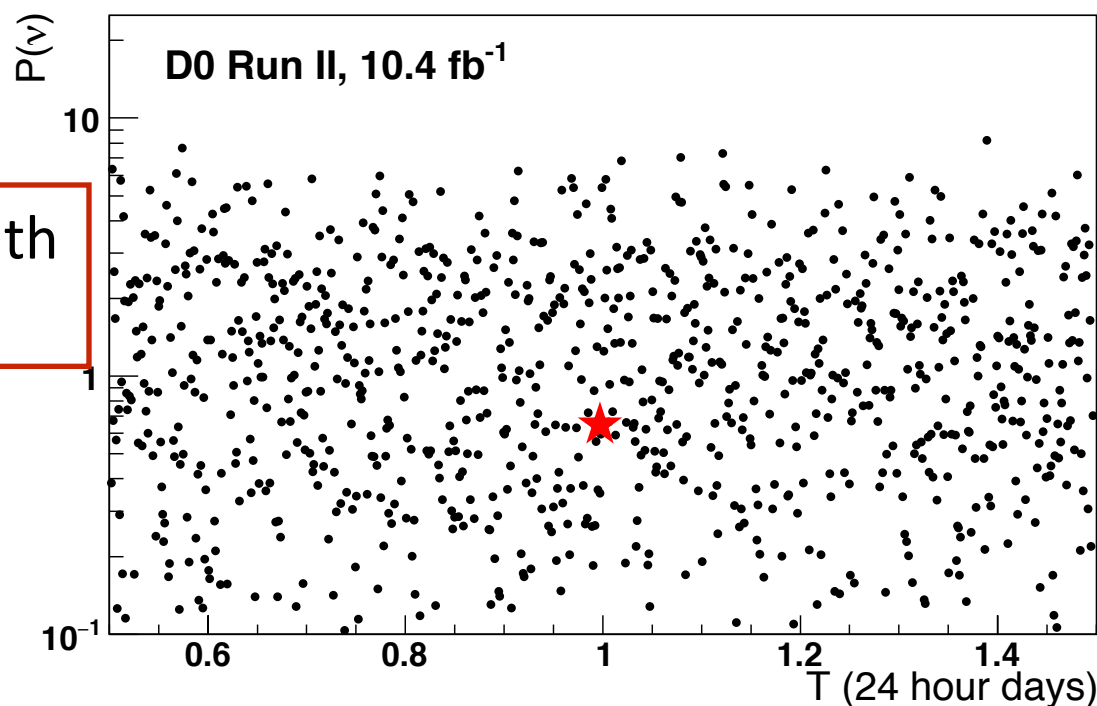
These limits imply that CPT violation is unlikely to contribute a significant fraction of the observed dimuon charge asymmetry, and that other explanations need to be sought.



# Periodogram

- Carry out cross check using periodogram  
- only good for varying part:

$$P(\nu) \equiv \frac{\left| \sum_{j=1}^N w_j \exp(-2\pi i \nu \hat{t}_j) \right|^2}{N \sigma_w^2},$$



$$\Delta a_{\perp} < 6.9 \times 10^{-13} \text{ GeV}$$



# Summary

- first limit placed on CPT violation exclusively in the  $B_s$  mixing system

$$\Delta a_{\perp} < 1.2 \times 10^{-12} \text{ GeV}$$

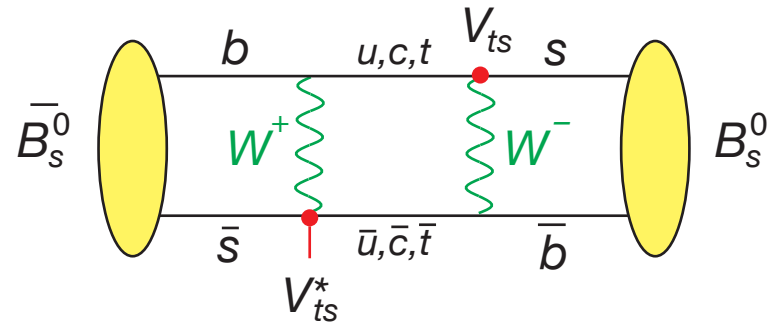
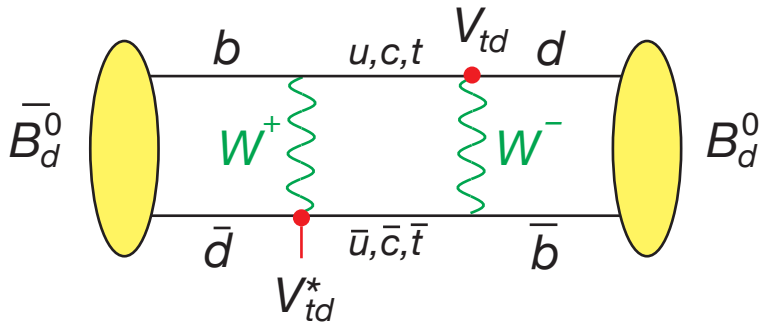
$$(-0.8 < \Delta a_T - 0.396 \Delta a_Z < 3.9) \times 10^{-13} \text{ GeV}$$

- No evidence of CPT violation
- Submitted to PRL: [arXiv:1506.04123 \[hep-ex\]](https://arxiv.org/abs/1506.04123)  
<http://www-d0.fnal.gov/Run2Physics/WWW/results/final/B/B15C/>

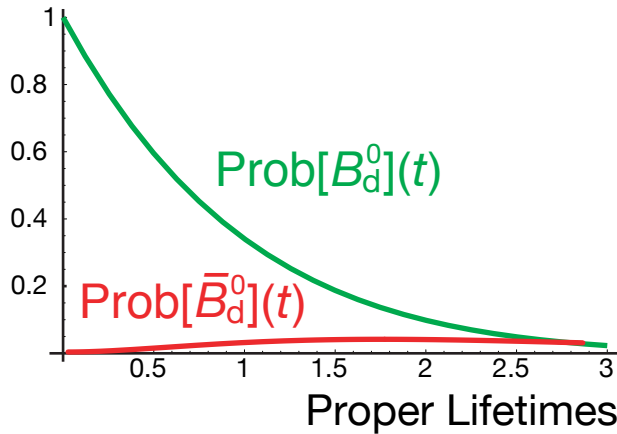


# Backup

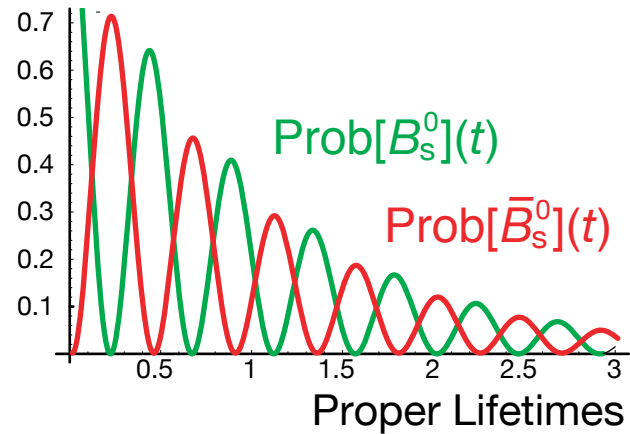
# $B^0\text{-}\bar{B}^0$ Mixing and Oscillations



- For  $B_d^0$



- For  $B_s^0$



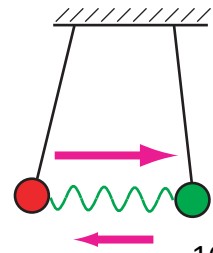
- With  $T$  (and therefore  $CP$ ) violation:

$$\left| \frac{p}{q} \right|^2 \neq 1,$$

$$|\langle M^0 | \bar{M}^0(t) \rangle|^2 \neq |\langle \bar{M}^0 | M^0(t) \rangle|^2$$

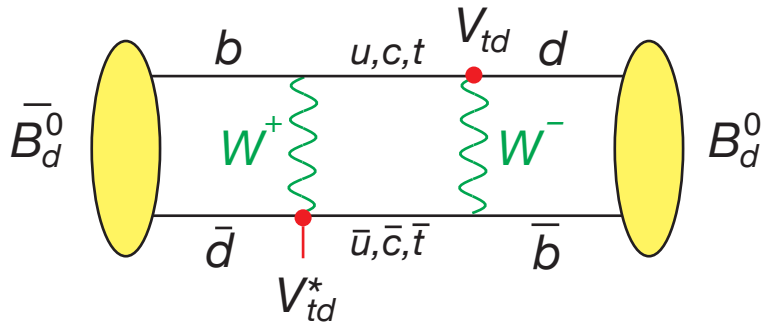
$$P(M^0 \rightarrow \bar{M}^0; t) \neq P(\bar{M}^0 \rightarrow M^0; t)$$

" $T$  ( $CP$ ) Violation in Mixing"

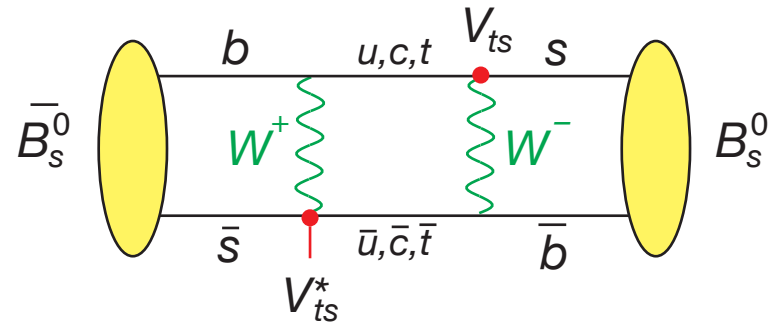
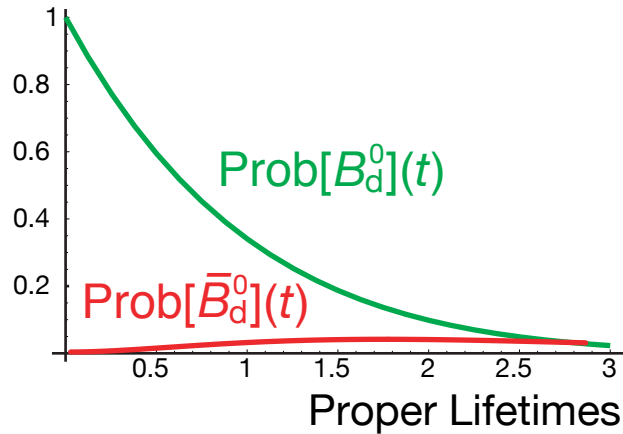




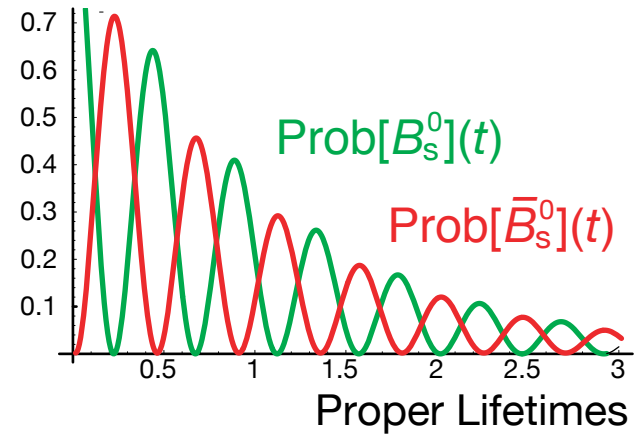
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$$P(M^0 \rightarrow M^0; t) \neq P(\bar{M}^0 \rightarrow \bar{M}^0; t)$$

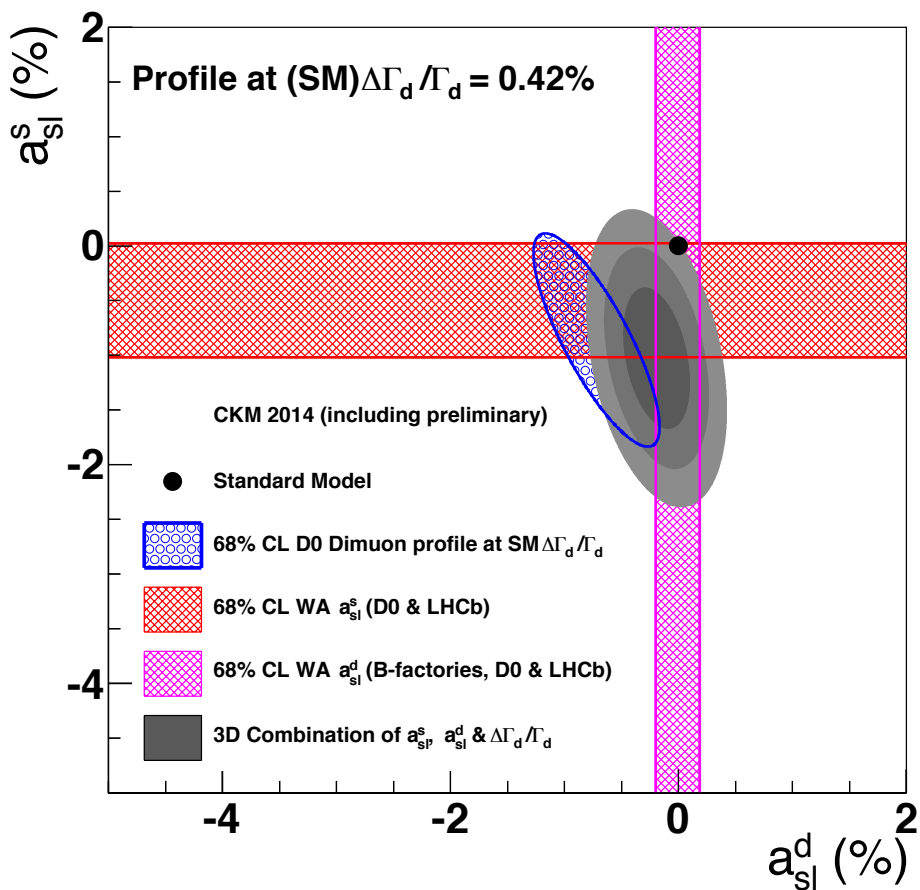
" $CPT$  Violation in Mixing"



# Introduction

- Interpretation of Charge asymmetries in single and like-sign dimuon events suggest that asymmetry due to  $B_s$  mesons.
- Currently no explanation!
- This could be explained by CPT violation.

V. Kostelecký and R Van Kooten  
Phys. Rev. D82: 101702, 2010



$$a_{sl}^s(\text{WA}) = (-0.83 \pm 0.43) \%,$$

$$a_{sl}^d(\text{WA}) = (-0.09 \pm 0.21) \%,$$

$$\Delta\Gamma_d/\Gamma_d(\text{WA}) = (1.51 \pm 0.91) \%,$$

$$\rho_{s,d} = -0.25,$$

$$\rho_{d,\Delta\Gamma} = +0.23$$

$$\rho_{s,\Delta\Gamma} = +0.47$$



# Parameters

TABLE I. Parameters and uncertainties in the extraction of the CPT-violating parameters.

Parameter	Value	Ref.
$A_0$	$(-0.40 \pm 0.31)\%$	Eq. 6
$A_1$	$(0.87 \pm 0.45)\%$	Eq. 6
$\phi$	$-2.28 \pm 0.51$	Eq. 6
$m_{B_s^0}$	$(5.36677 \pm 0.00024) \text{ GeV}$	[17]
$\Delta m_s$	$(17.761 \pm 0.022) \times 10^{12} \hbar s^{-1}$	[17]
$\Delta\Gamma_s/\Gamma_s$	$(0.138 \pm 0.012)$	[17]
$\hbar$	$6.58211928 \times 10^{-25} \text{ GeV}\cdot\text{s}$	[17]
$F_{B_s^0}^{\text{non-osc}} = F_{B_s^0}^{\text{osc}}$	$(0.465 \pm 0.017)$	[12]
$\langle p_z \rangle$	$(17.8 \pm 1.6) \text{ GeV}$	
$\langle p \rangle$	$(25.3 \pm 2.3) \text{ GeV}$	
Proton beam dir <sup>n</sup> $\alpha$	$219.53^\circ$	
Colatitude $\chi$	$48.17^\circ$	



# Systematic Uncertainties

- Fitting Uncertainties
- Mass Binning, Mass Range, Background polynomial, Fit Variations.

<b>Mass Range</b>	0.035%
<b>Mass Binning</b>	0.071%
<b>Fit Function</b>	0.085%
<b>Total</b>	<b>0.12%</b>

- Number of sidereal bins: relative uncertainty of central value of 8%
- Reconstruction asymmetries: relative uncertainty of 1%
- Added in quadrature